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Optical Focusing in Scattering Media with Photoacoustic Wavefront Shaping (PAWS)

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ABSTRACT

Controllable light delivery to the region of interest is essential to biomedical optical imaging methods like photoacoustic microscopy. It is, however, challenging beyond superficial depths in biological tissue (~1 mm beneath human skin) due to the strong scattering of light that scrambles the photon propagation paths. Recently, optical wavefront shaping has been proposed to modulate the incident light wavefront to compensate for the scattering-induced phase distortions, and consequentially, convey light optimally to a desired location behind or inside turbid media. To reach an optimum wavefront, a searching algorithm is usually required to optimize a feedback signal. In this work, we present our latest explorations, which use photoacoustic signals as the feedback to remotely and non-invasively guide the wavefront shaping process. Our method does not require direct optical access to the target region or the invasive embedding of fluorescence probes inside turbid media. Experimentally, we have demonstrated that diffuse light can be converged to the ultrasound focus by maximizing the amplitude of photoacoustic emissions from the intended absorbing site. Moreover, we show that wavefront-shaped light focusing can enhance existing optical imaging modalities like photoacoustic microscopy, in regard to signal-to-noise ratio, imaging depth, and potentially, resolution.

Keywords: optical focusing, wavefront shaping, photoacoustic imaging, light scattering, spatial light modulator, nonlinear photoacoustic signal, Grueneisen memory effect, optical speckle

1. INTRODUCTION

Light, in many ways, is an ideal choice to visualize biological structures, interrogate and control biological processes, and to diagnose and treat diseases. However, optical techniques lack spatial resolution in deep tissue due to the strong scattering of light from wavelength-scale refractive index changes in biological tissue [1]. As a result, when light is used alone, there is usually a trade-off between penetration depth and resolution; otherwise invasive procedures, such as embedded probes or guides, are required. Interdisciplinary methods have been developed to break this limitation, usually benefiting from the much weaker diffusivity of the mechanism other than light. Two types of ultrasound-mediated optical tomography—ultrasound-modulated optical tomography (UOT) [2] and photoacoustic imaging (PAI) [3, 4]—are examples. They have achieved optical contrast sensing with resolution provided by the externally applied ultrasound

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modulation [2] or the internally generated acoustic signal [3].

More recently, noticing that the appearance of random speckles formed by the propagation of coherent light in scattering media are actually deterministic within the speckle correlation time, researchers have started to explore the feasibility of optical focusing inside scattering media. For example, it has been shown that ultrasonically modulated/encoded light, as generated in UOT, can be time-reversed, and sent back to the scattering media, forming an optical focus at the ultrasonic focal position. This technique is usually called time-reversed ultrasonically encoded (TRUE) optical focusing [5-12], using either an analog (photorefractive material) or a digital (a digital camera plus a spatial light modulator) phase conjugation mirror to record and time-reverse the wavefront.

Another endeavor is optical wavefront shaping [13-23]. In this technique, the wavefront of an incident beam is shaped to compensate for the scattering-induced phase distortions, so that the scattered wavefronts interfere in phase at a predetermined location, and consequentially, form an optical focus. To reach the optimum wavefront compensation corresponding to a specific diffusivity pattern, a searching algorithm is usually employed to optimize a feedback signal. In this work, we present our preliminary explorations of using photoacoustic signals as the feedback to remotely and non-invasively guide the wavefront shaping process. This method, named photoacoustic wavefront shaping (PAWS), does not require either direct optical access to the target region or the invasive embedding of fluorescence probes inside scattering media. Experimentally, we show that diffuse light can be converged to the ultrasonic focus by maximizing the photoacoustic signal amplitude from the intended absorbing site.

2. EXPERIMENTAL SCHEMATIC AND PRINCIPLE

As shown in Fig. 1, the light from a pulsed laser source ($\lambda = 570$ nm, pulsed width = 5 ns, pulse energy ≈ 200 μ J, pulse repetition rate = 1 kHz) was split into two beams by a half-wave plate ($\lambda/2$ in Fig. 1) and a polarizing beam splitter (PBS). One beam, the vertically polarized portion reflected by the PBS, was detected by a photodiode to monitor the laser pulse energy fluctuations. The other beam, the horizontally polarized portion propagating through the PBS, was expanded and illuminated onto a spatial light modulator (SLM). The SLM has a full pixel resolution of 1920×1080 . In our experiments, the SLM was divided into 40×40 independently controlled blocks, each with linearized phase shifts between 0 and 2π . The reflected beam from the SLM was then focused onto a tissue-mimicking ground glass diffuser using a microscopic objective lens, creating a randomized speckle pattern behind the diffuser. A piece of black tape (part of the sample in Fig. 1) was positioned in water behind the diffuser, and served as the absorptive material to generate photoacoustic (PA) signals under the pulsed optical illumination. A focused ultrasound transducer (central frequency = 5 MHz, lateral focal width ≈ 600 μ m) was aligned to perceive PA emissions. The detected PA signals were amplified, digitized, and sent to a computer. The peak-to-peak amplitude of the acquired signals was quantified as the PA amplitudes, and was then normalized to the laser pulse energy. The measured amplitudes were then used as the feedback to guide the optimization of phase patterns on the SLM via a genetic algorithm [18, 22, 24]. Whenever a maximum PA amplitude was reached, the corresponding phase pattern on the SLM was considered the one that best compensated for the diffuser-induced scattering. Accordingly, an acoustic diffraction-limited optical focus at the transducer focal position would be expected. After the PAWS optimization, both the sample and the transducer were moved along the X direction to create a pathway for the CCD camera to image the light spots illuminating the semi-transparent tape section when different phase patterns were displayed on the SLM.

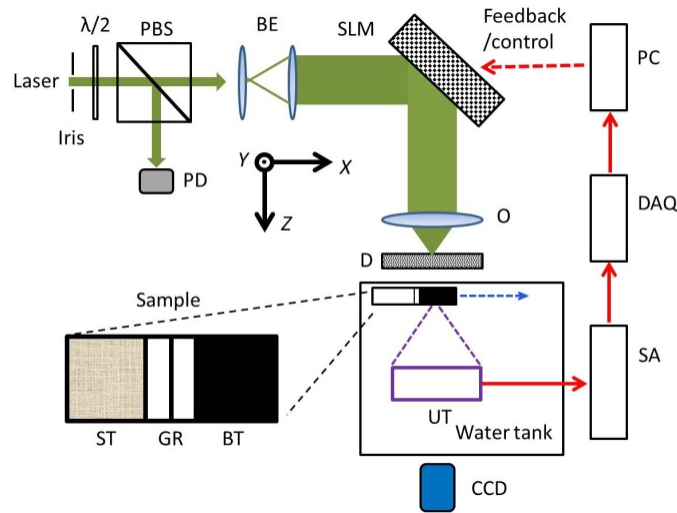


Fig. 1 Experimental schematic of PAWS. $\lambda/2$, half-wave plate; BE, beam expander; CCD, CCD camera; D, ground glass diffuser; DAQ, data acquisition device; O, microscopic objective lens; PBS, polarizing beam splitter; PC, computer; SA, photoacoustic signal amplifier; SLM, spatial light modulator; UT, focused ultrasonic transducer; XYZ , coordinate axes. The sample used to generate the photoacoustic signals during the optimization process is a piece of black tape (BT). After the PAWS, both the sample and the ultrasound transducer are moved along the X direction, so that the CCD camera is able to image the optical beam patterns illuminated on the semi-transparent tape (ST). Between the BT and ST is a graphite rod (GR, a pencil core 500 μm in diameter), which is used to calibrate the spatial scales of the beam pattern images.

3. EXPERIMENTAL VALIDATIONS

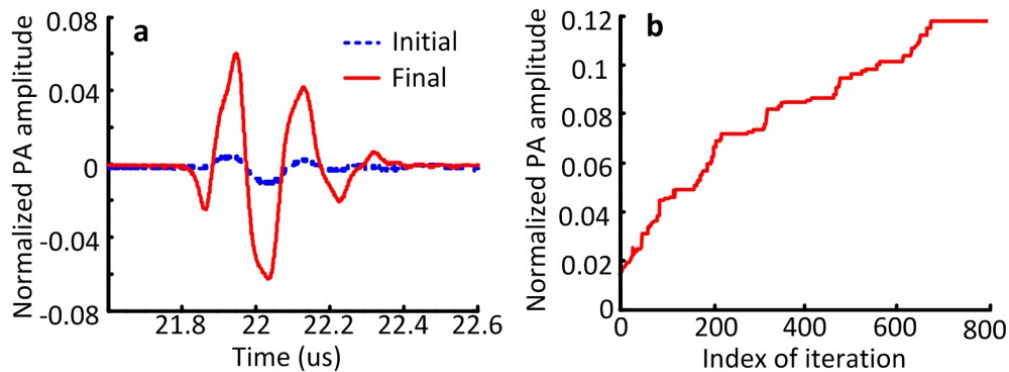


Fig. 2 (a) The initial and the final (optimized) PA signal before and after the PAWS optimization process. (b) Normalized PA amplitude versus index of iteration.

PAWS optical focusing was validated experimentally. First, a random phase pattern was displayed on the SLM. Accordingly, an initial PA signal (the blue, dashed curve in Fig. 2a) was recorded, after taking an average over 8 traces. As the genetic algorithm-based optimization proceeded, the normalized PA amplitude increased with the index of iteration (Fig. 2b). After ~ 700 iterations, the feedback plateaued. Therefore, the optimization was terminated after 800 iterations.

The PA signal at 800th iteration is shown in Fig. 2a. The signal enhancement factor between the final and the initial PA signal amplitude was $R \approx 0.12 / 0.014 \approx 8.6$, indicating that optical energy within the acoustical focus was enhanced by ~ 8.6 times with PAWS.

The enhancement factor was confirmed from the comparison of imaged patterns shown in Fig. 3. As seen, light was diffused and no optical focus could be seen (Fig. 3a) when the random phase pattern was displayed. But with the optimized phase pattern, an optical focus was clearly seen (Fig. 3b), with a full width at half-maximum (FWHM) of 610 μm and 602 μm along the X and Y directions, respectively (Fig. 3c-d), agreeing with the lateral focal width of the ultrasonic transducer (600 μm). This agreement demonstrates that the formed optical focus was acoustic-diffraction limited. Moreover, the peak-to-background ratio seen in Figs. 3b-d was around $1/0.12 \approx 8.3$, which was consistent with the 8.6 times enhancement in Fig. 2. It was also close to the theoretical enhancement factor estimated through the equation $\pi/4 \times (N/M)$ [7, 13], where $N = 20 \times 20$ is the number of independently controlled blocks on the SLM, and M is the number of speckle grains within the acoustic focus, which was ~ 119 in the current setting.

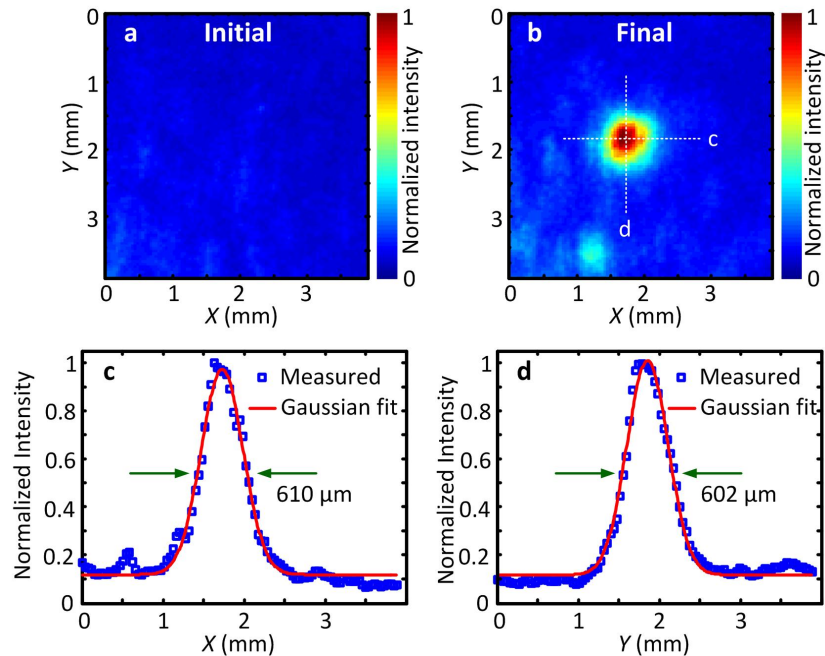


Fig. 3 The imaged beam pattern on the tape when a random (a) and the optimized phase pattern (b) were displayed on the SLM. Intensity distributions across the optical focus in (b) were shown in (c) and (d), along the X and Y axes respectively.

4. SUMMARY AND FUTURE WORK

We have shown in this work that one can use photoacoustic signals to guide optical wavefront shaping to achieve an acoustical diffraction-limited optical focus behind or inside scattering media. To break through this acoustic resolution limit, Conkey *et al.* [21] considered the non-uniform (Gaussian-like) spatial sensitivity profile of the ultrasonic transducer to discern and weight contributions to the final PA signals from each individual optical mode inside the ultrasonic focal region. Most recently, we proposed to use nonlinear photoacoustic signals, based on the Grueneisen memory effect, as

feedback for iterative optimization, which we termed nonlinear PAWS [25]. In our pilot experiment, an optically diffraction-limited optical focusing with superior peak fluence gain (~6000 times) was obtained in scattering media. Such an intense and highly confined optical focus in scattering media can benefit many micrometer-scale optical applications, especially if the optimization speed can be improved enough to enable biomedical applications in tissue, where optical speckles decorrelate fast due to physiological motions such as blood flow and aspiration.

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